# A Remark about Time-Analyticity of the Linear Landau Equation with Soft Potential 

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#### Abstract

In this note, we study the Cauchy problem of the linear spatially homogeneous Landau equation with soft potentials. We prove that the solution to the Cauchy problem enjoys the analytic regularizing effect of the time variable with an $L^{2}$ initial datum for positive time. So that the smoothing effect of Cauchy problem for the linear spatially homogeneous Landau equation with soft potentials is similar to the heat equation.


Key Words: Spatially homogeneous Landau equation, analytic smoothing effect, soft potentials. AMS Subject Classifications: 35B65, 76P05, 82C40

## 1 Introduction

The Cauchy problem of spatially homogenous Landau equation reads

$$
\left\{\begin{array}{l}
\partial_{t} F=Q(F, F),  \tag{1.1}\\
\left.F\right|_{t=0}=F_{0},
\end{array}\right.
$$

where $F=F(t, v) \geq 0$ is the density distribution function at time $t \geq 0$, with the velocity variable $v \in \mathbb{R}^{3}$. The Landau bilinear collision operator is defined by

$$
\begin{equation*}
Q(G, F)(v)=\sum_{j, k=1}^{3} \partial_{j}\left(\int_{\mathbb{R}^{3}} a_{j k}\left(v-v_{*}\right)\left[G\left(v_{*}\right) \partial_{k} F(v)-\partial_{k} G\left(v_{*}\right) F(v)\right] d v_{*}\right) \tag{1.2}
\end{equation*}
$$

with

$$
a_{j k}(v)=\left(\delta_{j k}|v|^{2}-v_{j} v_{k}\right)|v|^{\gamma}, \quad \gamma \geq-3,
$$

[^0]is a symmetric non-negative matrix such that
$$
\sum_{j, k=1}^{3} a_{j k}(v) v_{j} v_{k}=0
$$

Here, $\gamma$ is a parameter which leads to the classification of the hard potential if $\gamma>0$, Maxwellian molecules if $\gamma=0$, soft potential if $-3<\gamma<0$ and Coulombian potential if $\gamma=-3$.

The Landau equation was introduced as a limit of the Boltzmann equation when the collisions become grazing in $[6,18]$. In the hard potential case, the existence, and the uniqueness of the solution to the Cauchy problem for the spatially homogeneous Landau equation have been addressed by Desvillettes and Villani in [7,19]. Meanwhile, they also proved the smoothness of the solution is $C^{\infty}(] 0, \infty\left[; \mathcal{S}\left(\mathbb{R}^{3}\right)\right)$. The analytic and the Gevrey regularity of the solution for any $t>0$ have already been studied in [1,2].

We shall study the linearization of the Landau equation (1.1) near the Maxwellian distribution

$$
\mu(v)=(2 \pi)^{\frac{3}{2}} e^{-\frac{|v|^{2}}{2}} .
$$

Considering the fluctuation of the density distribution function

$$
F(t, v)=\mu(v)+\sqrt{\mu}(v) f(t, v),
$$

since $Q(\mu, \mu)=0$, the Cauchy problem (1.1) takes the form

$$
\left\{\begin{array}{l}
\partial_{t} f+\mathcal{L} f=\Gamma(f, f), \\
\left.f\right|_{t=0}=f_{0},
\end{array}\right.
$$

with $F_{0}=\mu+\sqrt{\mu} f_{0}$, where

$$
\begin{aligned}
& \Gamma(f, f)=\mu^{-\frac{1}{2}} Q(\sqrt{\mu} f, \sqrt{\mu} f) \\
& \mathcal{L}=\mathcal{L}_{1}+\mathcal{L}_{2} \quad \text { with } \mathcal{L}_{1} f=-\Gamma(\sqrt{\mu}, f), \quad \mathcal{L}_{2} f=-\Gamma(f, \sqrt{\mu}) .
\end{aligned}
$$

The spatially homogeneous Landau equation and non-cutoff Boltzmann equation in a close-to-equilibrium framework have been studied in [10] and the Gelfand-Shilov smoothing effect has been proved in [11,15]. Guo [8] constructed the classical solution for the spatially inhomogeneous Landau equation near a global Maxwellian in a periodic box. The smoothness of the solutions has been studied in [3,9,12]. In addition, the analytic smoothing effect of the velocity variable for the nonlinear Landau equation has been treated in $[13,16]$. The variant regularity results under a close-to-equilibrium setting have been considered in [4,5,17].

In this work, we consider the Cauchy problem of the linear Landau equation, such as

$$
\left\{\begin{array}{l}
\partial_{t} f+\mathcal{L} f=g,  \tag{1.3}\\
\left.f\right|_{t=0}=f_{0},
\end{array}\right.
$$

where $g$ is a analytic function with respect to the variable $t$ and $v$. The diffusion part $\mathcal{L}_{1}$ is written as follows

$$
\begin{equation*}
\mathcal{L}_{1} f=-\nabla_{v} \cdot\left[A(v) \nabla_{v} f\right]+\left(A(v) \frac{v}{2} \cdot \frac{v}{2}\right) f-\nabla_{v} \cdot\left[A(v) \frac{v}{2}\right] f \tag{1.4}
\end{equation*}
$$

with $A(v)=\left(\bar{a}_{j k}\right)_{1 \leq j, k \leq 3}$ is a symmetric matrix, and

$$
\bar{a}_{j k}=a_{j k} * \mu=\int_{\mathbb{R}^{3}}\left(\delta_{j k}\left|v-v^{\prime}\right|^{2}-\left(v_{j}-v_{j}^{\prime}\right)\left(v_{k}-v_{k}^{\prime}\right)\right)\left|v-v^{\prime}\right|^{\gamma} \mu\left(v^{\prime}\right) d v^{\prime} .
$$

We say that $u \in \mathcal{A}(\Omega)$ is an analytic function, where $\Omega \subset \mathbb{R}^{n}$ is an open domain, if $u \in C^{\infty}(\Omega)$ and there exists a constant $C$ such that for all multi-indices $\alpha \in \mathbb{N}^{n}$,

$$
\left\|\partial^{\alpha} u\right\|_{L^{\infty}(\Omega)} \leq C^{|\alpha|+1} \alpha!.
$$

Remark that, by using the Sobolev embedding, we can replace the $L^{\infty}$ norm by the $L^{2}$ norm, or norm in any Sobolev space in the above definition.

We study the linear Landau equation (1.3), with $-3<\gamma<0$, and show that the solution to the Cauchy problem (1.3) with the $L^{2}\left(\mathbb{R}^{3}\right)$ initial datum enjoys the analytic regularizing effect of the time variable. The main result reads as follows.
Theorem 1.1. For the soft potential $-3<\gamma<0$, for any $T>0$ and the initial datum $f_{0} \in$ $L^{2}\left(\mathbb{R}^{3}\right)$. Let $f$ be the solution of the Cauchy problem (1.3), then there exists a constant $C>0$ such that for any $k \in \mathbb{N}$, we have

$$
\begin{equation*}
\left\|\partial_{t}^{k} f(t)\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \leq \frac{C^{k+1}}{t^{k}} k!, \quad \forall t \in[0, T] \tag{1.5}
\end{equation*}
$$

For the linear operator with only the diffusion part of $\mathcal{L}_{1}$, the paper [14] prove that the Cauchy problem (1.3) admits a unique weak solution, and the solution satisfies for any $\alpha \in \mathbb{N}^{3}, \tilde{t}=\min (t, 1)$,

$$
\left\|\tilde{t}^{|\alpha|}\langle\cdot\rangle^{\frac{\gamma|\alpha|}{2}} \partial^{\alpha} f(t)\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \leq C^{|\alpha|+1} \alpha!, \quad \forall t>0 .
$$

With the similar computation, one can obtain the same analytical results as above, then using again the equation of (1.3), on have

$$
f \in C^{\infty}\left([0,+\infty] \mathcal{A}\left(\mathbb{R}^{3}\right)\right)
$$

So that we just need to prove (1.5) for the smooth solution of Cauchy problem (1.3).

## 2 Analysis of the Landau linear operator

In the following, the notation $A \lesssim B$ means there exists a constant $C>0$ such that $A \leq C B$. For simplicity, with $\gamma \in \mathbb{R}$, we denote the weighted Lebesgue spaces

$$
\left\|\langle\cdot\rangle^{\gamma} f\right\|_{L^{p}\left(\mathbb{R}^{3}\right)}=\|f\|_{p, \gamma}, \quad 1 \leq p \leq \infty,
$$

where we use the notation $\langle v\rangle=\left(1+|v|^{2}\right)^{\frac{1}{2}}$. And for the matrix $A$ defined in (1.4), we denote

$$
\|f\|_{A}^{2}=\sum_{j, k=1}^{3} \int\left(\bar{a}_{j k} \partial_{j} f \partial_{k} f+\frac{1}{4} \bar{a}_{j k} v_{j} v_{k} f^{2}\right) d v
$$

From corollary 1 in [8], for $\gamma>-3$, there exists a constant $C_{1}>0$ such that

$$
\begin{equation*}
\|f\|_{A}^{2} \geq C_{1}\left(\left\|\mathbf{P}_{v} \nabla f\right\|_{2, \frac{\gamma}{2}}^{2}+\left\|\left(\mathbf{I}-\mathbf{P}_{v}\right) \nabla f\right\|_{2,1+\frac{\gamma}{2}}^{2}+\|f\|_{2,1+\frac{\gamma}{2}}^{2}\right) \tag{2.1}
\end{equation*}
$$

where for any vector-valued function $G(v)=\left(G_{1}, G_{2}, G_{3}\right)$ define the projection to the vector $v=\left(v_{1}, v_{2}, v_{3}\right) \in \mathbb{R}^{3}$ as

$$
\left(\mathbf{P}_{v} G\right)_{j}=\sum_{k=1}^{3} G_{k} v_{k} \frac{v_{j}}{|v|^{2}}, \quad 1 \leq j \leq 3 .
$$

Since

$$
\nabla f=\mathbf{P}_{v} \nabla f+\left(\mathbf{I}-\mathbf{P}_{v}\right) \nabla f
$$

combining the inequality (2.1), we have

$$
\begin{equation*}
\|f\|_{A} \geq C_{1}\left(\|\nabla f\|_{2, \frac{\gamma}{2}}+\|f\|_{2,1+\frac{\gamma}{2}}\right) \tag{2.2}
\end{equation*}
$$

For later use, we need the following results for the coefficients to the linear Landau operator, which have been proved in [14].
Lemma 2.1 ([14]). For any $\beta \in \mathbb{R}^{3}$ with $|\beta| \geq 1$ and $\bar{a}_{j k}$ was defined in (1.4) with $-3<\gamma<0$, then we have

$$
\begin{equation*}
\left|\partial^{\beta} \bar{a}_{j k}(v)\right| \lesssim\langle v\rangle^{\gamma+1} \sqrt{\beta!} . \tag{2.3}
\end{equation*}
$$

Moreover, for any $\beta \in \mathbb{R}^{3}$,

$$
\begin{align*}
& \left|\partial^{\beta}\left(\sum_{j, k=1}^{3} \partial_{j} a_{j k} *\left(v_{k} \mu\right)\right)\right| \lesssim\langle v\rangle^{\gamma+1}(|\beta|+1) \sqrt{\beta!},  \tag{2.4a}\\
& \left|\partial^{\beta}\left(\sum_{j, k=1}^{3} \bar{a}_{j k} v_{j} v_{k}\right)\right| \lesssim\langle v\rangle^{\gamma+1}(|\beta|+1) \sqrt{\beta!} . \tag{2.4b}
\end{align*}
$$

Lemma 2.2 ([14]). Let $f_{1}, f_{2} \in \mathcal{S}\left(\mathbb{R}^{3}\right)$, $\bar{a}_{j k}$ was defined in (1.4) with $-3<\gamma<0$. For any $\beta \in \mathbb{R}^{3}$, we have

$$
\begin{equation*}
\left|\sum_{j, k=1}^{3}\left(\partial^{\beta} \bar{a}_{j k} \partial_{k} f_{1}, \partial_{j} f_{2}\right)_{L^{2}\left(\mathbb{R}^{3}\right)}\right| \lesssim \sqrt{\beta!}\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{A} \tag{2.5}
\end{equation*}
$$

By using the results of the coefficients to the linear Landau operator in [14], we can obtain the following estimates. Firstly, for any $\gamma>-3$ and $\delta>0$, we have

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}|v-w|^{\gamma} e^{-\delta|w|^{2}} d w \lesssim\langle v\rangle^{\gamma} . \tag{2.6}
\end{equation*}
$$

Lemma 2.3. Let $f \in \mathcal{S}\left(\mathbb{R}^{3}\right)$, and $-3<\gamma<0$, then for any $0<\epsilon_{1}<1$, there exists a constant $C_{\epsilon_{1}}>0$ such that

$$
\left(1-\epsilon_{1}\right)\|f\|_{A}^{2} \leq\left(\mathcal{L}_{1} f, f\right)_{L^{2}}+C_{\epsilon_{1}}\|f\|_{2, \frac{\gamma}{2}}^{2} .
$$

Proof. By using the representation (1.4), and integrating by parts, we have

$$
\begin{aligned}
-\left(\mathcal{L}_{1} f, f\right)_{L^{2}} & =-\int_{\mathbb{R}^{3}}\left(\bar{a}_{j k} \partial_{j} f \partial_{k} f+\frac{1}{4} \bar{a}_{j k} v_{j} v_{k} f^{2}\right)-\frac{1}{2} \int_{\mathbb{R}^{3}} \partial_{j}\left(\bar{a}_{j k} v_{k}\right) f^{2} \\
& =-\|f\|_{A}^{2}+R_{0} .
\end{aligned}
$$

Since

$$
\begin{equation*}
\sum_{j} a_{j k} v_{j}=\sum_{k} a_{j k} v_{k}=0 \tag{2.7}
\end{equation*}
$$

we have

$$
\partial_{j}\left(\bar{a}_{j k} v_{k}\right)=\partial_{j}\left(a_{j k} *\left(v_{k} \mu\right)\right)=\partial_{j} a_{j k} *\left(v_{k} \mu\right)
$$

Therefore from (2.4) and the Cauchy-Schwarz inequality, it follows that

$$
\left|R_{0}\right| \lesssim \int_{\mathbb{R}^{3}}\langle v\rangle^{\gamma+1} f^{2}(v) d v \lesssim\|f\|_{2, \frac{\gamma}{2}}\|f\|_{2,1+\frac{\gamma}{2}},
$$

then by using (2.2) and the Cauchy-Schwarz inequality, for any $0<\epsilon_{1}<1$, we have

$$
\left|R_{0}\right| \leq C_{2}\|f\|_{2, \frac{\gamma}{2}}\|f\|_{A} \leq \epsilon_{1}\|f\|_{A}^{2}+\frac{4 C_{2}^{2}}{\epsilon_{1}}\|f\|_{2, \frac{\gamma}{2}}^{2} .
$$

Let $C_{\epsilon_{1}}=\frac{4 C_{2}^{2}}{\epsilon_{1}}$, then we can conclude

$$
\left(1-\epsilon_{1}\right)\|f\|_{A}^{2} \leq\left(\mathcal{L}_{1} f, f\right)_{L^{2}}+C_{\epsilon_{1}}\|f\|_{2, \frac{\gamma}{2}}^{2} .
$$

This completes the proof.
Proposition 2.1. Let $f_{1}, f_{2} \in \mathcal{S}\left(\mathbb{R}^{3}\right)$ and $-3<\gamma<0$, then there exists a constant $C_{2}>0$ such that

$$
\left|\left(\mathcal{L}_{1} f_{1}, f_{2}\right)_{L^{2}}\right| \leq C_{2}\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{A} .
$$

Proof. By using the representation (1.4), and integrating by parts, we have

$$
\begin{aligned}
\left(\mathcal{L}_{1} f_{1}, f_{2}\right)_{L^{2}}= & \int_{\mathbb{R}^{3}} \bar{a}_{j k} \partial_{j} f_{1} \partial_{k} f_{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} \bar{a}_{j k} v_{j} v_{k} f_{1} f_{2} \\
& +\frac{1}{2} \int_{\mathbb{R}^{3}} \partial_{j}\left(\bar{a}_{j k} v_{k}\right) f_{1} f_{2} \\
= & R_{1}+R_{2}+R_{3} .
\end{aligned}
$$

Since $-3<\gamma<0$, by the inequality (2.5), we obtain

$$
\left|R_{1}\right| \lesssim\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{A}
$$

For the term $R_{2}$ and $R_{3}$, using (2.7), then from (2.3) and (2.4), it follows that

$$
\left|R_{2}\right|+\left|R_{3}\right| \lesssim \int_{\mathbb{R}^{3}}\langle v\rangle^{\gamma+1}\left|f_{1}(v) f_{2}(v)\right| d v,
$$

then using the Cauchy-Schwarz inequality and (2.2), we have

$$
\left|R_{2}\right|+\left|R_{3}\right| \lesssim\left\|f_{1}\right\|_{2,1+\frac{\gamma}{2}}\left\|f_{2}\right\|_{2,1+\frac{\gamma}{2}} \lesssim\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{A} .
$$

Finally, combining $R_{1}-R_{3}$ to get

$$
\left|\left(\mathcal{L}_{1} f_{1}, f_{2}\right)_{L^{2}}\right| \leq C_{2}\left\|f_{1}\right\|\left\|_{A}\right\| f_{2} \|_{A} .
$$

This completes the proof.
Now, we shall estimate $\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}$. Firstly, we give the representation of the operator $\mathcal{L}_{2}$. For $f \in \mathcal{S}\left(\mathbb{R}^{3}\right)$, from (1.2), it follows that

$$
\begin{align*}
\mathcal{L}_{2} f & =-\mu^{-\frac{1}{2}} Q(\sqrt{\mu} f, \mu) \\
& =\mu^{-\frac{1}{2}} \partial_{j}\left(\int_{\mathbb{R}^{3}} a_{j k}\left(v-v^{\prime}\right)\left[\mu^{\frac{1}{2}}\left(v^{\prime}\right) f\left(v^{\prime}\right) v_{k}+\partial_{k}\left(\mu^{\frac{1}{2}} f\right)\left(v^{\prime}\right)\right] d v^{\prime} \mu(v)\right) \\
& =\mu^{-\frac{1}{2}} \partial_{j}\left[\mu\left(a_{j k} *\left(v_{k} \mu^{\frac{1}{2}} f\right)+\partial_{k} a_{j k} *\left(\mu^{\frac{1}{2}} f\right)\right)\right] . \tag{2.8}
\end{align*}
$$

Proposition 2.2. Let $f_{1}, f_{2} \in \mathcal{S}\left(\mathbb{R}^{3}\right)$ and $-3<\gamma<0$, then there exists a constant $C_{3}>0$ such that

$$
\begin{equation*}
\left|\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}\right| \leq C_{3}\left(\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{A}+\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}}\right) \tag{2.9}
\end{equation*}
$$

Proof. Using integration by parts with (2.8), we have

$$
\begin{aligned}
\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}= & -\left(a_{j k} *\left(v_{k} \mu^{\frac{1}{2}} f_{1}\right), \mu^{\frac{1}{2}}\left(\frac{v_{j}}{2} f_{2}+\partial_{j} f_{2}\right)\right)_{L^{2}} \\
& +\left(\partial_{j k} a_{j k} *\left(\mu^{\frac{1}{2}} f_{1}\right), \mu^{\frac{1}{2}} f_{2}\right)_{L^{2}}-\left(\partial_{k} a_{j k} *\left(\mu^{\frac{1}{2}} f_{1}\right), \mu^{\frac{1}{2}} v_{j} f_{2}\right)_{L^{2}} \\
= & I_{1}
\end{aligned}+I_{2}+I_{3} .
$$

Since

$$
\begin{equation*}
\left|\partial^{\alpha} a_{j k}(v)\right| \leq|v|^{\gamma+2-|\alpha|}, \quad \forall \alpha \in \mathbb{N}^{3}, \tag{2.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\langle v\rangle^{\beta} \mu^{\rho}(v) \in L^{\infty}\left(\mathbb{R}^{3}\right), \quad \forall \beta \in \mathbb{R}, \quad \rho>0, \tag{2.11}
\end{equation*}
$$

by Cauchy-Schwarz inequality, we obtain

$$
\begin{aligned}
\left|a_{j k} *\left(v_{k} \mu^{\frac{1}{2}} f_{1}\right)\right| & \leq \int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\gamma+2}\left\langle v^{\prime}\right\rangle^{1-\frac{\gamma}{2}} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left\langle v^{\prime}\right\rangle^{\frac{\gamma}{2}}\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& \lesssim\left(\int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{2(\gamma+2)} \mu^{\frac{1}{2}}\left(v^{\prime}\right) d v^{\prime}\right)^{\frac{1}{2}}\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} .
\end{aligned}
$$

For any $-3<\gamma<0$, we have $2(\gamma+2)>-3$, then by using the inequality (2.6), it follows that

$$
\left|a_{j k} *\left(v_{k} \mu^{\frac{1}{2}} f_{1}\right)\right| \lesssim\langle v\rangle^{\gamma+2}\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} .
$$

Using the Cauchy-Schwarz inequality and the inequality (2.2), we can conclude

$$
\begin{aligned}
\left|I_{1}\right| & \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} \int_{\mathbb{R}^{3}}\langle v\rangle^{\gamma+3} \mu^{\frac{1}{2}}(v)\left(\left|f_{2}(v)\right|+\left|\nabla f_{2}(v)\right|\right) d v \\
& \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left(\left\|f_{2}\right\|_{2,1+\frac{\gamma}{2}}+\left\|\nabla f_{2}\right\|_{2, \frac{\gamma}{2}}\right) \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{A} .
\end{aligned}
$$

For the term $I_{2}$, from (2.10), one has

$$
\left|\partial_{j k} a_{j k} *\left(\mu^{\frac{1}{2}} f_{1}\right)\right| \lesssim \int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\gamma} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} .
$$

Consider two sets $\left\{\left|v-v^{\prime}\right| \leq 1\right\}$ and $\left\{\left|v-v^{\prime}\right| \geq 1\right\}$, that is

$$
\begin{aligned}
& \int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\gamma} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
= & \int_{\left|v-v^{\prime}\right| \leq 1}+\int_{\left|v-v^{\prime}\right| \geq 1}=A_{1}+A_{2} .
\end{aligned}
$$

For the term $A_{1}$, since $-3<\gamma<0$, we have

$$
\begin{aligned}
A_{1} & =\sum_{j \leq 0} \int_{2^{j-1} \leq\left|v-v^{\prime}\right| \leq 2^{j}}\left|v-v^{\prime}\right|^{\gamma} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& \leq \sum_{j \leq 0}\left(2^{j-1}\right)^{\gamma} \int_{\left|v-v^{\prime}\right| \leq 2^{j}} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& =8 \sum_{j \leq 0}\left(2^{j-1}\right)^{\gamma+3} \frac{1}{2^{3 j}} \int_{\left|v-v^{\prime}\right| \leq 2^{j}} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& \leq 8 \sum_{j \leq 0}\left(2^{j-1}\right)^{\gamma+3} M\left(\mu^{\frac{1}{2}} f_{1}\right) \lesssim M\left(\mu^{\frac{1}{2}} f_{1}\right),
\end{aligned}
$$

where $M$ is the Hardy-Littlewood maximal function. For term $A_{2}$, from (2.6),

$$
\begin{aligned}
A_{2} & =\int_{\left|v-v^{\prime}\right| \geq 1}\left|v-v^{\prime}\right|^{\gamma} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& \leq \int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\gamma+2} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} \\
& \lesssim\left(\int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{2(\gamma+2)} \mu^{\frac{1}{2}}\left(v^{\prime}\right) d v^{\prime}\right)^{\frac{1}{2}}\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} \\
& \lesssim\langle v\rangle^{\gamma+2}\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} .
\end{aligned}
$$

Combining $A_{1}$ and $A_{2}$, using Cauchy-Schwarz inequality to get

$$
\begin{aligned}
\left|I_{2}\right| & \lesssim \int_{\mathbb{R}^{3}}\left|M\left(\mu^{\frac{1}{2}} f_{1}\right)(v) \mu^{\frac{1}{2}}(v) f_{2}(v)\right| d v+\left\|f_{1}\right\|_{2, \frac{\gamma}{2}} \int_{\mathbb{R}^{3}}\langle v\rangle^{\gamma+2} \mu^{\frac{1}{2}}(v)\left|f_{2}(v)\right| d v \\
& \lesssim\left\|M\left(\mu^{\frac{1}{2}} f_{1}\right)\right\|_{L^{2}}\left\|\mu^{\frac{1}{2}} f_{2}\right\|_{L^{2}}+\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}} \\
& \lesssim\left\|\mu^{\frac{1}{2}} f_{1}\right\|_{L^{2}}\left\|\mu^{\frac{1}{2}} f_{2}\right\|_{L^{2}}+\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}} \\
& \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}} \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{A} .
\end{aligned}
$$

For $I_{3}$, from (2.10), we have

$$
\left|\partial_{k} a_{j k} *\left(\mu^{\frac{1}{2}} f_{1}\right)\right| \lesssim \int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\gamma+1} \mu^{\frac{1}{2}}\left(v^{\prime}\right)\left|f_{1}\left(v^{\prime}\right)\right| d v^{\prime} .
$$

Note that

$$
\frac{3}{2}(\gamma+1)>-3
$$

with $-3<\gamma<0$. Using Hölder's inequality, (2.11) and (2.6), we have

$$
\begin{aligned}
\left|\partial_{k} a_{j k} *\left(\mu^{\frac{1}{2}} f_{1}\right)\right| & \lesssim\left(\int_{\mathbb{R}^{3}}\left|v-v^{\prime}\right|^{\frac{3}{2}(\gamma+1)} \mu^{\frac{1}{2}}\left(v^{\prime}\right) d v^{\prime}\right)^{\frac{2}{3}}\left\|f_{1}\right\|_{3, \frac{\gamma}{2}} \\
& \lesssim\langle v\rangle^{\gamma+1}\left\|f_{1}\right\|_{3, \frac{\gamma}{2}} .
\end{aligned}
$$

Now, we want to show $\left\|f_{1}\right\|_{3, \frac{\gamma}{2}}$ can be bounded by $\left\|f_{1}\right\|_{A}$. By applying Hölder's inequality, $\langle v\rangle^{\gamma / 2} f_{1}(v)$ in $L^{3}\left(\mathbb{R}^{3}\right)$ can be bounded by

$$
\left(\left\|\langle\cdot\rangle^{\gamma / 2} f_{1}\right\|_{L^{2}}\left\|\langle\cdot\rangle^{\gamma / 2} f_{1}\right\|_{L^{6}}\right)^{\frac{1}{2}}
$$

and Sobolev embedding implies

$$
\left\|\langle\cdot\rangle^{\gamma / 2} f_{1}\right\|_{L^{6}} \lesssim\left\|\nabla\left[\langle\cdot\rangle^{\gamma / 2} f_{1}\right]\right\|_{L^{2}}
$$

thus we get

$$
\left\|f_{1}\right\|_{3, \frac{\gamma}{2}} \lesssim\left(\left\|\langle\cdot\rangle^{\gamma / 2} f_{1}\right\|_{L^{2}}\left\|\nabla\left[\langle\cdot\rangle^{\gamma / 2} f_{1}\right]\right\|_{L^{2}}\right)^{\frac{1}{2}}
$$

Notice that

$$
\nabla\left[\langle v\rangle^{\gamma / 2} f_{1}(v)\right]=\frac{\gamma}{2}\langle v\rangle^{\gamma / 2-2} f_{1}(v) v+\langle v\rangle^{\gamma / 2} \nabla f_{1}(v),
$$

from (2.2), we have

$$
\left\|\nabla\left[\langle\cdot\rangle^{\gamma / 2} f_{1}\right]\right\|_{L^{2}} \leq\left\|\frac{\gamma}{2}\langle v\rangle^{\gamma / 2-2} f_{1} v\right\|_{L^{2}}+\left\|\langle\cdot\rangle^{\gamma / 2} \nabla f_{1}\right\|_{L^{2}} \lesssim\left\|f_{1}\right\|_{A},
$$

which implies

$$
\begin{equation*}
\left\|f_{1}\right\|_{3, \frac{\gamma}{2}} \lesssim\left\|f_{1}\right\|_{A} . \tag{2.12}
\end{equation*}
$$

Finally, using the Cauchy-Schwarz inequality and (2.12) to get

$$
\left|I_{3}\right| \lesssim\left\|f_{1}\right\|_{A} \int_{\mathbb{R}^{3}}\langle v\rangle^{\gamma+2} \mu^{\frac{1}{2}}(v)\left|f_{2}(v)\right| d v \lesssim\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}} .
$$

Combining $I_{1}-I_{3}$, we obtain

$$
\left|\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}\right| \leq C_{3}\left(\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{A}+\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{2, \frac{\gamma}{2}}\right) .
$$

This completes the proof.

## Remark 2.1.

(1). For $f_{1}, f_{2} \in \mathcal{S}\left(\mathbb{R}^{3}\right)$ and $\gamma>-5 / 2$, we have

$$
\left|\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}\right| \lesssim\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}\left\|f_{2}\right\|_{A} .
$$

(2). For $-3<\gamma<0$, if $f_{1}=f_{2}$, then for any $\epsilon_{2}>0$, there exists a constant $C_{\epsilon_{2}}>0$ such that

$$
\begin{equation*}
\left|\left(\mathcal{L}_{2} f_{1}, f_{1}\right)_{L^{2}}\right| \leq \epsilon_{2}\left\|f_{1}\right\|_{A}^{2}+C_{\epsilon_{2}}\left\|f_{1}\right\|_{2, \frac{\gamma}{2}}^{2} . \tag{2.13}
\end{equation*}
$$

(3). From (2.2),

$$
\begin{equation*}
\left|\left(\mathcal{L}_{2} f_{1}, f_{2}\right)_{L^{2}}\right| \leq C_{4}\left\|f_{1}\right\|_{A}\left\|f_{2}\right\|_{A} . \tag{2.14}
\end{equation*}
$$

## 3 Energy estimates

In this section, we study the energy estimates of the solution to the Cauchy problem (1.3).
Lemma 3.1. For $-3<\gamma<0$. Lef $f$ be the solution of Cauchy problem (1.3). Assume $f_{0} \in$ $L^{2}\left(\mathbb{R}^{3}\right)$. Then there exists a constant $C_{5}>0$ such that for any $T>0$ and $t \in[0, T]$,

$$
\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\int_{0}^{t}\|f(s)\|_{A}^{2} d s \leq\left(C_{5}\right)^{2} .
$$

Proof. Since $f$ is the solution of Cauchy problem (1.3),

$$
\frac{1}{2} \frac{d}{d t}\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\left(\mathcal{L}_{1} f, f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}=(g, f)_{L^{2}\left(\mathbb{R}^{3}\right)}-\left(\mathcal{L}_{2} f, f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}
$$

Since $\gamma<0$, by using Lemma 2.3 and (2.13), we have

$$
\begin{aligned}
& \quad \frac{d}{d t}\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+2\left(1-\epsilon_{1}\right)\|f(t)\|_{A}^{2} \\
& \leq 2\|f(t)\|_{L^{2}}\|g(t)\|_{L^{2}}+2 \epsilon_{2}\|f(t)\|_{A}^{2}+2\left(C_{\epsilon_{1}}+C_{\epsilon_{2}}\right)\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}
\end{aligned}
$$

Applying Cauchy-Schwarz inequality, and choosing $\epsilon_{1}=\epsilon_{2}=\frac{1}{4}$, we can get

$$
\frac{d}{d t}\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\|f(t)\|_{A}^{2} \leq c_{1}\|f(t)\|_{L^{2}}^{2}+\|g(t)\|_{L^{2}}^{2}
$$

Since $g$ is analytic with respect to $t$ and $v$, for any $T>0$ and $t \in[0, T]$, there exists a constant $A>0$ such that

$$
\begin{equation*}
\|g(t)\|_{L^{2}} \leq A \tag{3.1}
\end{equation*}
$$

Therefore by applying Gronwall inequality, for any $T>0$ and $t \in[0, T]$, taking

$$
C_{5} \geq e^{\frac{1}{2} c_{1} T} \sqrt{\left\|f_{0}\right\|_{L^{2}}^{2}+T A^{2}}
$$

one can obtain

$$
\|f(t)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\int_{0}^{t}\|f(s)\|_{A}^{2} d s \leq e^{c_{1} T}\left(\left\|f_{0}\right\|_{L^{2}}^{2}+T A^{2}\right) \leq\left(C_{5}\right)^{2}
$$

This completes the proof.
Lemma 3.2. For $-3<\gamma<0$. Lef $f$ be the solution of Cauchy problem (1.3). Assume $f_{0} \in$ $L^{2}\left(\mathbb{R}^{3}\right)$. Then there exists a constant $C_{6}>0$ such that for any $T>0$ and $t \in[0, T]$,

$$
\begin{equation*}
\left\|t \partial_{t} f\right\|_{\left.L^{\infty}(j 0, T] ; L^{2}\left(\mathbb{R}^{3}\right)\right)}^{2}+\int_{0}^{T}\left\|t \partial_{t} f\right\|_{A}^{2} d t \leq\left(C_{6}\right)^{2} \tag{3.2}
\end{equation*}
$$

Proof. Since the solution of Cauchy problem (1.3) belongs to $C^{\infty}(] 0, T\left[; \mathcal{S}\left(\mathbb{R}^{3}\right)\right)$, we have that

$$
\partial_{t}\left(t \partial_{t} f\right)+\mathcal{L}_{1}\left(t \partial_{t} f\right)=\partial_{t} f-\mathcal{L}_{2}\left(t \partial_{t} f\right)+t \partial_{t} g,
$$

and for $0<t \leq T$,

$$
\begin{aligned}
& \frac{1}{2}\left\|t \partial_{t} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\int_{0}^{t}\left(\mathcal{L}_{1}\left(s \partial_{s} f\right), s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
= & \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s-\int_{0}^{t}\left(\mathcal{L}_{2}\left(s \partial_{s} f\right), s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}+\int_{0}^{t}\left(s \partial_{s} f, s \partial_{s} g\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
= & S_{1}+S_{2}+S_{3} .
\end{aligned}
$$

Firstly, since $\gamma<0$, by Lemma 2.3 for all $0<t \leq T$, we can conclude

$$
\begin{aligned}
& \int_{0}^{t}\left(\mathcal{L}_{1}\left(s \partial_{s} f\right), s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\geq & \left(1-\epsilon_{1}\right) \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s-C_{\epsilon_{1}} \int_{0}^{t}\left\|s \partial_{s} f\right\|_{2, \frac{\gamma}{2}}^{2} d s \\
\geq & \left(1-\epsilon_{1}\right) \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s-T C_{\epsilon_{1}} \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s .
\end{aligned}
$$

For the term $S_{1}$, since $f$ is the solution of (1.3), using Proposition 2.1 and (2.14), for all $0<t \leq T$, we have

$$
\begin{aligned}
& \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \\
= & \int_{0}^{t}\left(g, s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s-\int_{0}^{t}\left(\mathcal{L}_{1} f, s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s-\int_{0}^{t}\left(\mathcal{L}_{2} f, s \partial_{s} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\leq & \int_{0}^{t}\|g(s)\|_{L^{2}\left(\mathbb{R}^{3}\right)}\left\|s \partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} d s+C_{2} \int_{0}^{t}\|f(s)\|_{A}\left\|s \partial_{s} f\right\|_{A} d s \\
\quad & +C_{4} \int_{0}^{t}\|f(s)\|_{A}\left\|s \partial_{s} f\right\|_{A} d s .
\end{aligned}
$$

For all $0<t \leq T$, by Cauchy-Schwarz inequality, it follows that

$$
\begin{aligned}
& \int_{0}^{t}\|g(s)\|_{L^{2}\left(\mathbb{R}^{3}\right)}\left\|s \partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\leq & T^{\frac{1}{2}} \int_{0}^{t}\|g(s)\|_{L^{2}\left(\mathbb{R}^{3}\right)} s^{\frac{1}{2}}\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\leq & \frac{1}{2} \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+\frac{T}{2} \int_{0}^{t}\|g(s)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s .
\end{aligned}
$$

Using Cauchy-Schwarz inequality, since $\gamma<0$, for any $0<\delta<1$, we have

$$
\begin{aligned}
& \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \\
\leq & \delta \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s+T \int_{0}^{t}\|g(s)\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+C_{\delta} \int_{0}^{t}\|f(s)\|_{A}^{2} d s,
\end{aligned}
$$

with $C_{\delta}$ depends on $C_{2}, C_{4}$. Combining (3.1) and Lemma 3.1, for $0<\delta<1$,

$$
\begin{equation*}
S_{1}=\int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \leq \delta \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s+T^{2} A^{2}+C_{\delta} C_{5}^{2} \tag{3.3}
\end{equation*}
$$

For the term $S_{2}$, let $f_{1}=s \partial_{s} f$ in (2.13), then for all $0<t \leq T$,

$$
\left|S_{2}\right| \leq \epsilon_{2} \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s+C_{\epsilon_{2}} T \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s
$$

by using (3.3) with $c_{3} T \delta \leq \epsilon_{2}$,

$$
\left|S_{2}\right| \leq 2 \epsilon_{2} \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s+\tilde{C}_{\epsilon_{2}}
$$

with $\tilde{C}_{\varepsilon_{2}}$ depends on $C_{2}, C_{4}, C_{5}, A$ and $T$.
Finally, for the term $S_{3}$, by Cauchy-Schwarz inequality, it follows that

$$
\left|S_{3}\right| \leq \int_{0}^{t}\left\|s \partial_{s} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+T \int_{0}^{t} s\left\|\partial_{s} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s .
$$

Since $g$ is analytic with respect to $t$ and $v$, for all $0<t \leq T$, we have

$$
\left\|t \partial_{t} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \leq A^{2}
$$

applying (3.3) with $T \delta \leq \epsilon_{2}$ to get $S_{3}$ can be bounded by

$$
\epsilon_{2} \int_{0}^{t}\left\|s \partial_{s} f\right\|_{A}^{2} d s+T A^{4}+T\left(T^{2} A^{2}+C_{\epsilon_{2}} C_{5}^{2}\right)
$$

Therefore, combining the results above and using (3.3) with $T C_{\epsilon_{1}} \delta<\epsilon_{1}$, let $\epsilon_{1}=\epsilon_{2}=\frac{1}{16}$, $0<\delta \leq \frac{1}{4}$, and taking $C_{6} \geq \sqrt{\tilde{C}_{5}}$, we get

$$
\left\|t \partial_{t} f\right\|_{\left.L^{\infty}(j 0, T] ; L^{2}\left(\mathbb{R}^{3}\right)\right)}^{2}+\int_{0}^{T}\left\|t \partial_{t} f\right\|_{A}^{2} d t \leq \tilde{C}_{5} \leq\left(C_{6}\right)^{2} .
$$

with $C_{6}$ depend on $C_{2}, C_{4}, C_{5}, A$ and $T$.

## 4 Analytic smoothing effect for time variable

In this section, we will show the analytic regularity of the time variable for $t>0$. We construct the following estimate, which implies Theorem 1.1 immediately.
Proposition 4.1. For $-3<\gamma<0$. Let $f$ be the solution of Cauchy problem (1.3), and $f_{0} \in$ $L^{2}\left(\mathbb{R}^{3}\right)$. Then there exists a constant $B>0$ such that for any $T>0, t \in[0, T]$ and $k \in \mathbb{N}$,

$$
\begin{equation*}
\left\|t^{k} \partial_{t}^{k} f\right\|_{\left.L^{\infty}(j 0, T] ; L^{2}\left(\mathbb{R}^{3}\right)\right)}^{2}+\int_{0}^{T}\left\|t^{k} \partial_{t}^{k} f\right\|_{A}^{2} d t \leq\left(B^{k+1} k!\right)^{2} \tag{4.1}
\end{equation*}
$$

Proof. We prove this proposition by induction on the index $k$. For $k=1$, it is enough to take in (3.2). Assume (4.1) holds true, for any $1 \leq m \leq k-1$ with $k \geq 2$,

$$
\begin{equation*}
\left\|t^{m} \partial_{t}^{m} f\right\|_{L^{\infty}\left((0, T] ; L^{2}\left(\mathbb{R}^{3}\right)\right)}^{2}+\int_{0}^{T}\left\|t^{m} \partial_{t}^{m} f\right\|_{A}^{2} d t \leq\left(B^{m+1} m!\right)^{2} \tag{4.2}
\end{equation*}
$$

We shall prove (4.1) holds true for $m=k$.

Since $\mu$ is the function with respect to $v$, which implies

$$
t^{k} \partial_{t}^{k} \mathcal{L}_{1} f=\mathcal{L}_{1}\left(t^{k} \partial_{t}^{k} f\right), \quad t^{k} \partial_{t}^{k} \mathcal{L}_{2} f=\mathcal{L}_{2}\left(t^{k} \partial_{t}^{k} f\right)
$$

Then by (1.3), we have

$$
\partial_{t}\left(t^{k} \partial_{t}^{k} f\right)+\mathcal{L}_{1}\left(t^{k} \partial_{t}^{k} f\right)=k t^{k-1} \partial_{t}^{k} f-\mathcal{L}_{2}\left(t^{k} \partial_{t}^{k} f\right)+t^{k} \partial_{t}^{k} g .
$$

Taking the $L^{2}\left(\mathbb{R}^{3}\right)$ inner product of both sides with respect to $t^{k} \partial_{t}^{k} f$, we get

$$
\begin{aligned}
& \frac{1}{2} \frac{d}{d t}\left\|t^{k} \partial_{t}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\left(\mathcal{L}_{1}\left(t^{k} \partial_{t}^{k} f\right), t^{k} \partial_{t}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} \\
= & k\left(t^{k-1} \partial_{t}^{k} f, t^{k} \partial_{t}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}-\left(\mathcal{L}_{2}\left(t^{k} \partial_{t}^{k} f\right), t^{k} \partial_{t}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}+\left(t^{k} \partial_{t}^{k} g, t^{k} \partial_{t}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} .
\end{aligned}
$$

For all $0<t \leq T$, integrating from 0 to $t$, since $\gamma<0$, by using Lemma 2.3, it follows that

$$
\begin{aligned}
& \int_{0}^{t}\left(\mathcal{L}_{1}\left(s^{k} \partial_{s}^{k} f\right), s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\geq & \left(1-\epsilon_{1}\right) \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s-C_{\epsilon_{1}} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{2, \frac{\gamma}{2}}^{2} d s \\
\geq & \left(1-\epsilon_{1}\right) \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s-T C_{\epsilon_{1}} \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s,
\end{aligned}
$$

and let $f_{1}=s^{k} \partial_{s}^{k} f$ in (2.13) to get

$$
\int_{0}^{t}\left|\left(\mathcal{L}_{2}\left(s^{k} \partial_{s}^{k} f\right), s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}\right| d s \leq \epsilon_{2} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s+T C_{\epsilon_{2}} \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s
$$

then using Cauchy-Schwarz inequality to get

$$
\begin{aligned}
& \int_{0}^{t}\left|\left(s^{k} \partial_{s}^{k} g, s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)}\right| d s \\
\leq & \frac{1}{2} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+\frac{1}{2} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \\
\leq & \frac{1}{2} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+\frac{T}{2} \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s .
\end{aligned}
$$

Combining the results above, and taking $\epsilon_{1}=\epsilon_{2}=\frac{1}{8}$, we have for all $0<t \leq T$,

$$
\begin{align*}
& \left\|t^{k} \partial_{t}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\frac{3}{2} \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s \\
\leq & \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+C_{7} \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s, \tag{4.3}
\end{align*}
$$

with $C_{7}$ depends on $T$.

Since $f$ is the solution of (1.3) and $k \geq 2$, we have

$$
\partial_{t}^{k} f=\partial_{t}^{k-1} g-\mathcal{L}_{1}\left(\partial_{t}^{k-1} f\right)-\mathcal{L}_{2}\left(\partial_{t}^{k-1} f\right),
$$

which implies

$$
\begin{aligned}
& \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \\
= & \int_{0}^{t}\left(s^{k-1} \partial_{s}^{k-1} g, s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s-\int_{0}^{t}\left(\mathcal{L}_{1}\left(s^{k-1} \partial_{s}^{k-1} f\right), s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
& \quad-\int_{0}^{t}\left(\mathcal{L}_{2}\left(s^{k-1} \partial_{s}^{k-1} f\right), s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s .
\end{aligned}
$$

For all $0<t \leq T$, using Cauchy-Schwarz inequality, we have

$$
\begin{aligned}
& \int_{0}^{t}\left(s^{k-1} \partial_{s}^{k-1} g, s^{k} \partial_{s}^{k} f\right)_{L^{2}\left(\mathbb{R}^{3}\right)} d s \\
\leq & T^{\frac{1}{2}} \int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \frac{s k}{2 k-1}_{2}^{l}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \\
\leq & \frac{T}{2} \int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\frac{1}{2} \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s .
\end{aligned}
$$

By using Proposition 2.1, (2.14) and Cauchy-Schwarz inequality, for any $0<\delta<1$, there exists a constant $C_{\delta}>0$ such that for all $0<t \leq T$,

$$
\begin{align*}
& \int_{0}^{t} s^{2 k-1}\left\|\partial_{s}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s \\
\leq & \delta \int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s+T \int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+C_{\delta} \int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} f\right\|_{A}^{2} d s, \tag{4.4}
\end{align*}
$$

with $C_{\delta}$ depends on $C_{2}, C_{4}$. Let $C_{7} \delta \leq \frac{1}{2}$, substituting (4.4) into (4.3), we get

$$
\begin{aligned}
& \left\|t^{k} \partial_{t}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s \\
\leq & \tilde{C}_{7}\left(\int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s+\int_{0}^{t}\left\|s^{k-1} \partial_{s}^{k-1} f\right\|_{A}^{2} d s\right)+\int_{0}^{t}\left\|s^{k} \partial_{s}^{k} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2} d s,
\end{aligned}
$$

with $\tilde{C}_{7}$ depends on $C_{2}, C_{4}, C_{7}$ and $T$.
Finally, since $g$ is analytic with respect to $t$ and $v$, for any $k \in \mathbb{N}$, there exists a constant $A>0$ such that for any $0<t \leq T$,

$$
\left\|t^{k} \partial_{t}^{k} g\right\|_{L^{2}\left(\mathbb{R}^{3}\right)} \leq A^{k+1} k!,
$$

taking $B \geq \max \left\{A, \sqrt{2 \tilde{C}_{7}}\right\}$, using the induction hypothesis (4.2), we obtain

$$
\begin{aligned}
& \left\|t^{k} \partial_{t}^{k} f\right\|_{L^{2}\left(\mathbb{R}^{3}\right)}^{2}+\int_{0}^{t}\left\|s^{k} \partial_{s}^{k} f\right\|_{A}^{2} d s \\
\leq & \tilde{C}_{7}\left(\left(A^{k}(k-1)!\right)^{2}+\left(B^{k}(k-1)!\right)^{2}\right)+\left(A^{k+1} k!\right)^{2} \\
\leq & \left(B^{k+1} k!\right)^{2},
\end{aligned}
$$

with $B$ depends on $C_{1}, C_{2}, C_{4}, A$ and $T$. We finish the proof of Proposition 4.1.

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